

# OH stretch IR spectra of (H<sub>2</sub>O)<sub>3</sub> and benzene-(H<sub>2</sub>O)<sub>3</sub>

S.Y. Fredericks<sup>1</sup>, J.M. Pedulla<sup>1</sup>, K.D. Jordan<sup>1</sup>, T.S. Zwier<sup>2</sup>

<sup>1</sup>Department of Chemistry, University of Pittsburgh, Pittsburgh, PA 15260, USA

<sup>2</sup>Department of Chemistry, Purdue University, West Lafayette, IN 47907, USA

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**Abstract.** The OH stretch IR spectrum of (H<sub>2</sub>O)<sub>3</sub> is significantly altered upon complexation of the cluster with a benzene molecule. Theoretical calculations show that the differences in the OH stretch spectra of (H<sub>2</sub>O)<sub>3</sub> and benzene-(H<sub>2</sub>O)<sub>3</sub> are primarily due to mode localization induced by the benzene molecule.

**Key words:** Clusters – Hydrogen bonding – Vibrational frequencies

## 1 Introduction

Recent years have witnessed a sharp upsurge in research on small water clusters. This has been driven by experimental advances, permitting characterization of the clusters in the microwave [1–4], near IR [5, 6], and far IR [7–10] regions, as well as by computational advances, enabling accurate ab initio calculations [11–22] of the structures, stabilities, and vibrational spectra of the clusters.

Mixed clusters containing water and organic molecules are also receiving increased attention. Benzene-(H<sub>2</sub>O)<sub>n</sub> (hereafter referred to as BW<sub>n</sub>) clusters, in particular, have been the subject of a number of experimental [23–32] and theoretical studies [33–36]. Near IR spectra have been published for clusters as large as BW<sub>7</sub> [30–32]. The hydrogen bonding topologies of the W<sub>n</sub> cluster sub-units in the BW<sub>n</sub> clusters closely mimic those of the pure W<sub>n</sub> clusters of the same size. Since the OH stretch infrared spectra are sensitive to the number, type, and strength of the hydrogen bonds present, they serve as sensitive spectral signatures for various fundamental H-bonding topologies. The BW<sub>n</sub> clusters themselves provide examples of H-bonded cycles, cubes, and three-dimensional networks.

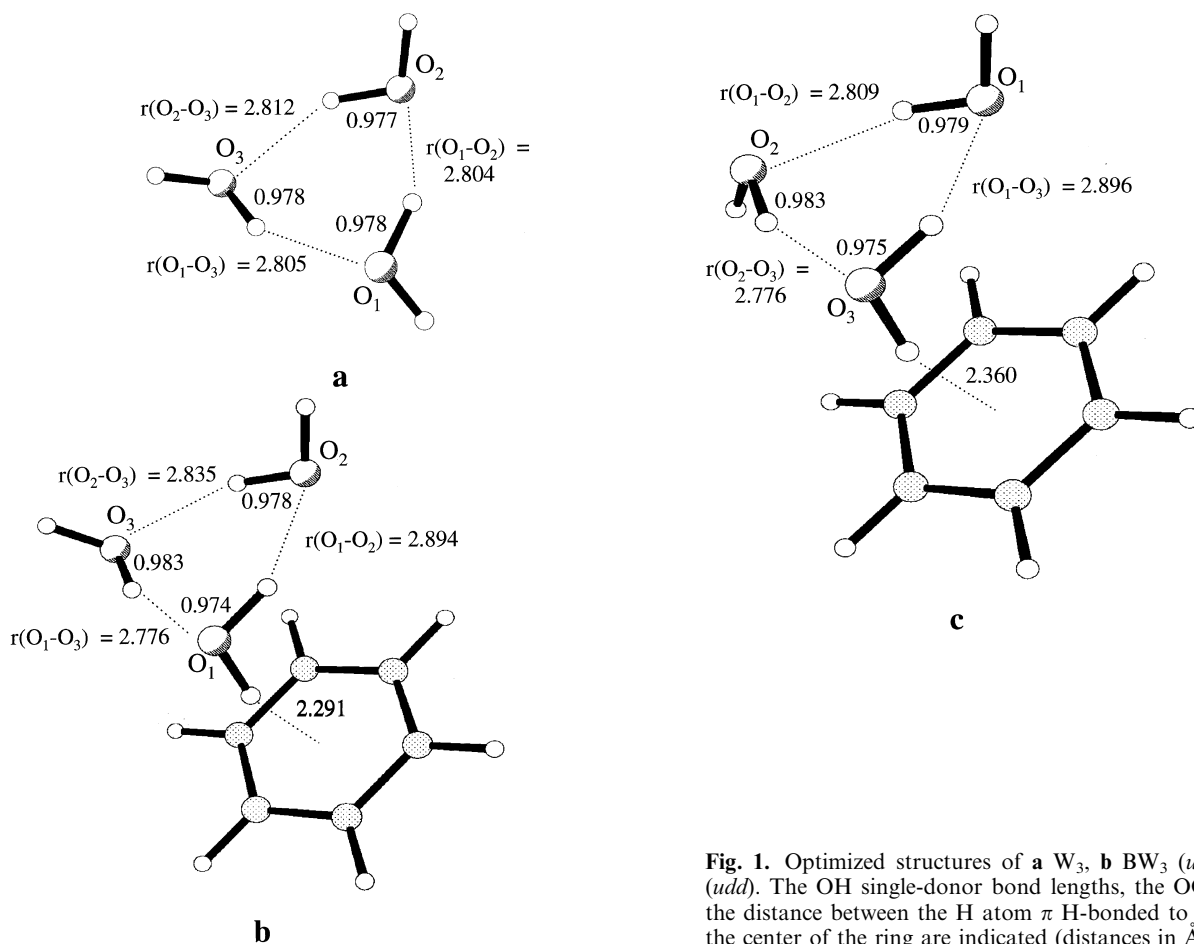
One important question which arises in the study of these fundamental H-bonding topologies is the sensitivity of their spectral signatures to their local environment. In particular, one wonders whether the presence of a solute molecule will perturb them, and if so, in what way. Given the tremendous interest in W<sub>n</sub> clusters and the availability of data on the BW<sub>n</sub> clusters, the comparison of ab initio results on the two systems can provide just such insight into the perturbing effects of benzene on W<sub>n</sub> clusters. Furthermore, such studies of the mixed clusters provide insight into how interactions at a molecular level influence macroscopic behavior such as miscibility. Finally, the availability of spectroscopic and high-quality ab initio data on the pure and mixed clusters will prove valuable in the development of new, improved model potentials for carrying out Monte Carlo and molecular dynamics simulations.

In every case studied so far, comparison of the OH stretch IR data for the BW<sub>n</sub> and W<sub>n</sub> clusters indicates that the lowest energy structure for the W<sub>n</sub> cluster is retained in the presence of benzene [30–32]. This is consistent with the observation that the interaction between water molecules is greater than that between benzene and water molecules (the estimated D<sub>0</sub> values of the W<sub>2</sub> and BW<sub>1</sub> dimers are 3.6 and 2.2 kcal/mol, respectively [37, 38]). Nonetheless, the presence of the benzene molecule can induce significant changes in the OH stretch spectra of the water clusters [30–32, 35]. Intuitively, we expect that the impact of the benzene molecule on the OH stretch spectra will be greater for highly symmetrical than for non-symmetrical water clusters. This follows from the observation that in the high symmetry clusters, the OH stretch normal modes are highly delocalized over all water molecules in the cluster, while addition of the benzene will cause partial mode localization. In contrast, for the non-symmetrical clusters, the OH stretch vibrations are already localized, and the interaction with the benzene should cause relatively small perturbations on the spectra. This is borne out experimentally [30–32] as well as by calculations [35].

The simplest water cluster for which complexation with benzene could potentially cause significant mode localization is the dimer [30, 31, 34]. The dimer possesses

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Correspondence to: K.D. Jordan



**Fig. 1.** Optimized structures of **a**  $W_3$ , **b**  $BW_3$  (*udd*), and **c**  $BW_3$  (*udd*). The OH single-donor bond lengths, the OO distances, and the distance between the H atom  $\pi$  H-bonded to the benzene and the center of the ring are indicated (distances in Å)

$C_s$  symmetry, and as a result, the two OH groups of the acceptor monomer are symmetrically equivalent. In the global minimum structure of  $BW_2$  the  $W_2$  molecule  $\pi$  H-bonds to the benzene ring through one of the H atoms of the acceptor group, and this breaks the symmetry. In the absence of tunneling between different permutational isomers, the two OH stretch modes on the acceptor monomer would then be localized. However, the barriers for tunneling are sufficiently low [36] for it to be likely that under experimental conditions [30, 31] these modes remain delocalized. The dimer with symmetrically equivalent OH groups on one monomer is a special case; more common are clusters in which two or more water monomers are symmetrically equivalent or nearly so. Members of this class of clusters include the cyclic water trimer, tetramer, and pentamer, and the cubic forms of the octamer.

In the present investigation,  $W_3$  is chosen as representative of water clusters for which the complexation with a benzene molecule causes a large perturbation of the vibrational spectrum.  $W_3$  has a cyclic global minimum structure, with each monomer acting as a single H-donor as shown in Fig. 1. Two of the free OH groups are on one side of the OOO plane, and the third is on the other side [9–12, 21, 22, 39–41]; as a result, the global minimum structure is of  $C_1$  symmetry. Nonetheless, the cluster (even in a single potential minimum) may be viewed as having approximate threefold symmetry, as

evidenced by the fact that the three OO distances are nearly equal, as are the three single-donor OH bond lengths [11, 12, 39]. Moreover, the barriers for flipping the H atoms of the free OH groups and for torsional exchange of the free and single-donor OH groups of the monomers are low, making possible rapid interconversion of the equivalent global minima [12, 21, 22, 39, 40]. The molecular symmetry group appropriate for the tunneling spectra is isomorphic with the  $C_{3h}$  point group [9].

In Ref. [35], we reported theoretical structures and normal-mode OH stretch spectra of  $W_n$  and  $BW_n$ ,  $n = 1-3$ , clusters. The present study extends this earlier work by examining more closely the localization of the OH stretching modes in  $BW_3$ , by presenting the calculated spectrum of two different forms of  $BW_3$ , and by reporting also the calculated vibrational spectra in regions other than the OH stretch regions. Visualization of the normal modes is particularly useful for understanding the changes induced by the benzene, and animations of the normal modes are accessible via the Internet archive.

## 2 Structural information

All theoretical results reported here were obtained from MP2 calculations employing the 6-31+G[2d,p] basis set. The details of the calculations are presented in Ref.

[35]. Figure 1 depicts the structures, including the OO and OH distances, of  $W_3$  and of two isomers of  $BW_3$ , labeled as  $BW_3(uud)$  and  $BW_3(udd)$ . The “*u*” and “*d*” denote, respectively, whether a particular free OH group is pointing up or down with respect to the OOO plane, where the direction of *u* and *d* is set by benzene being placed below the  $W_3$  cycle [41]. For the free  $W_3$  species the (*uud*) and (*udd*) species are two of the 96 energetically equivalent minima in the relevant subgroup of the full permutation inversion group [39–41]. However, in both  $BW_3$  isomers, the  $W_3$  entity is  $\pi$  H-bonded to the benzene ring through one of the free OH groups, thereby breaking the degeneracy between  $BW_3(uud)$  and  $BW_3(udd)$ . The  $BW_3(uud)$  isomer is the global minimum on the  $BW_3$  potential energy surface, whereas the  $BW_3(udd)$  isomer is calculated to lie  $217\text{ cm}^{-1}$  higher in energy. It follows from this that the tunneling dynamics for  $BW_3$  will be very different from that of the bare  $W_3$  cluster. In fact, the most important tunneling motion in  $BW_3$  most likely involves motion of the  $W_3$  relative to the benzene rather than motion of the H atoms of the  $W_3$  portion of the cluster.

As expected, the interaction with the benzene ring destroys the approximate threefold symmetry present in the bare  $W_3$  cluster. For example, the three OO distances in the  $BW_3(uud)$  cluster are 2.776, 2.835, and 2.894 Å, whereas those in the bare  $W_3$  cluster fall in the narrow range of 2.804 to 2.812 Å. Complete structural information is available electronically via the Internet archive.

### 3 Vibrational spectra

Figure 2 compares the OH stretch spectra of  $W_3$  and the (*uud*) isomer of  $BW_3$ . The figure includes both the MP2 [35] and experimental results [30, 31]. The spectral transitions are reported relative to the mean of the two OH stretch frequencies (calculated or experimental as appropriate) of water monomer. The complete calculated spectra for  $W_3$ ,  $BW_3(uud)$  and  $BW_3(udd)$  are available via the Internet archive.

A detailed analysis of the OH stretch spectra of  $W_3$  and  $BW_3$  has been reported in Ref. [35], and only the most important details are repeated here. The IR spectrum of  $W_3$  is characterized by two pronounced transitions in the OH stretch region (with shifts of  $+14$  and  $-173\text{ cm}^{-1}$ ). Based on the calculations, it is seen that one of these is due to the three free OH stretch modes, which are nearly degenerate in the MP2 calculations (shifts  $+14$  to  $+20\text{ cm}^{-1}$ ). The second intense transition is due to two nearly degenerate single-donor OH stretch modes (calculated shifts:  $-237$ ,  $-249\text{ cm}^{-1}$ ). The third single-donor OH stretch mode, which involves an in-phase oscillation of the three OH groups, is even further shifted than the other two single-donor modes. It carries near zero intensity (as a result of the approximate threefold symmetry of the cluster) and is not observed in the experimental studies.

The OH stretch spectrum (both calculated and measured) of  $BW_3$  is markedly different from that of  $W_3$ . The experimental spectrum shows five intense lines, the assignment of which is straightforward. The band due to

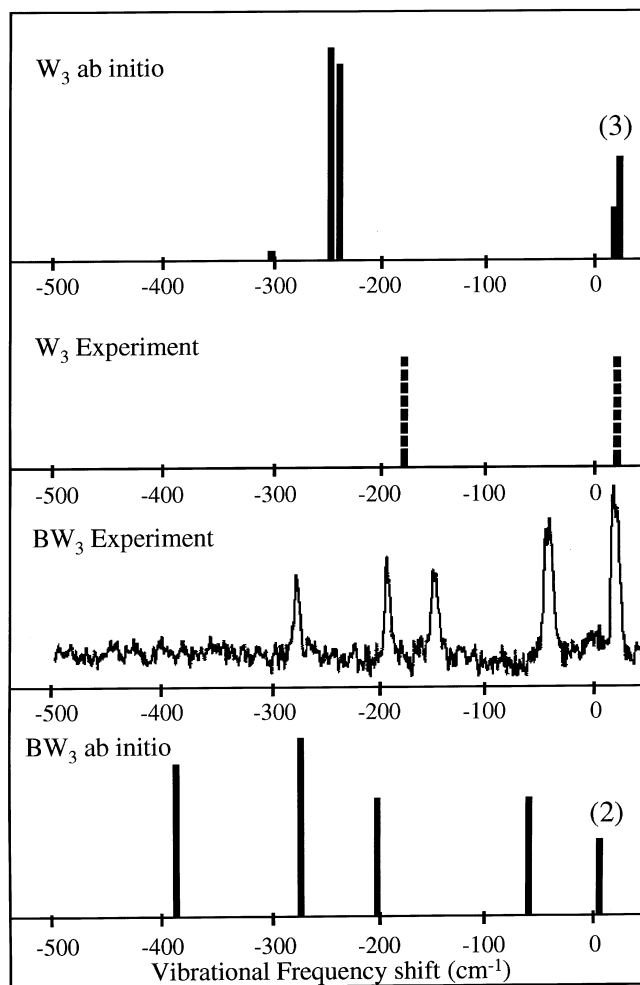


Fig. 2. Calculated (MP2 level) and experimental vibrational frequency shifts for the OH stretch region of  $W_3$  and  $BW_3(uud)$ . The shifts are with respect to the average of the OH stretch frequencies of the water monomer, from theory ( $3900\text{ cm}^{-1}$ ) or experiment ( $3706\text{ cm}^{-1}$ ), as appropriate. Experimental results are from Ref. [1] for  $W_1$ , Ref. [6] for  $W_3$ , and Ref. [32] for  $BW_3$ .

the three free OH stretch transitions of  $W_3$  is split into two bands near  $+10$  and  $-49\text{ cm}^{-1}$  in  $BW_3$ . The latter “split-off” transition is associated with the OH group  $\pi$  H-bonded to the benzene ring. The presence of the benzene also causes the three single-donor OH stretch transitions to spread out and to acquire nearly equal intensities.

The MP2 calculations closely reproduce the changes in the OH stretch spectra in going from  $W_3$  to  $BW_3$ . The one shortcoming of the MP2 calculations is that they predict larger shifts (in the negative direction) in the single-donor OH stretch frequencies than are observed experimentally. This is due to the underestimation of the OO distances at this level of theory; this deficiency can be remedied by the inclusion of higher-order correlation effects [35].

The (*uud*) and (*udd*) forms of  $BW_3$  have similar OH stretch spectra. The most important difference is in the free OH stretch region of the spectrum. In the (*uud*) isomer two of the free OH stretch modes remain nearly

degenerate, whereas in the (*udd*) isomer the splittings among the three free OH stretch lines are sufficiently large to allow the observation of three distinct bands in the experimental spectrum (if this isomer were present in sufficiently high amount under the experimental conditions). This difference between the (*uud*) and (*udd*) isomers is a consequence of the fact that there are two free OH groups pointing toward the benzene ring in the latter isomer but only one, that  $\pi$ -hydrogen bonded to the benzene, pointed toward the ring in the former isomer.

The MP2 calculations predict that the free OH stretch modes are localized in both  $W_3$  and  $BW_3$ . In the case of the bare  $W_3$  cluster, this is a consequence of the very weak coupling between the three free OH groups and the fact that the MP2 calculations neglect tunneling between the equivalent potential energy minima. The free OH stretch modes of  $W_3$  would be found to be delocalized if averaging over the low-frequency tunneling motions were included. In contrast, the localization of the free OH modes of  $BW_3$  would “survive” averaging over the low-frequency motions. The greatest effect of the interaction with the benzene molecule is on the OH stretch mode associated with the OH group  $\pi$  H-bonded to the benzene ring. This mode is shifted to lower frequency and experiences a large increase in intensity. The large intensity change derives from variation of either the polarization of the electron distribution of the benzene or the extent of charge transfer between the water cluster and the benzene molecule accompanying motion along this OH stretch coordinate [35].

The MP2 calculations predict the single-donor OH stretch modes to be delocalized in the bare  $W_3$  molecule, but to be highly localized in the (*uud*) and (*udd*)  $BW_3$  isomers. In the mixed clusters over 80% of the weight of each mode is associated with a single OH group. MPEG animations of the six OH stretch modes of  $W_3$  and of the global minimum structure of  $BW_3$  are available electronically via the Internet archive.

The splitting pattern for the three single-donor OH stretch modes of the bare  $W_3$  can be understood by consideration of the simplified Hückel-like Hamiltonian matrix

$$\mathbf{H} = \begin{pmatrix} \alpha & \beta & \beta \\ \beta & \alpha & \beta \\ \beta & \beta & \alpha \end{pmatrix}, \quad (1)$$

where the diagonal ( $\alpha$ ) entries refer to the frequencies of the OH stretch modes localized on each of the three monomers, and the off-diagonal ( $\beta$ ) entries describe the coupling between the localized OH stretch modes. The three diagonal entries were set equal as were the various off-diagonal entries to account for the effective threefold symmetry of the cluster. Comparison of the eigenvalues ( $\alpha - \beta$ ,  $\alpha - \beta$ ,  $\alpha + 2\beta$ ) of this Hamiltonian with the calculated spectra gives a value of  $23 \text{ cm}^{-1}$  for the coupling. This is about half as large as the coupling ( $50 \text{ cm}^{-1}$ ) between the two localized OH stretch modes of the water monomer.

The overall splitting among the three single-donor OH stretch modes of  $BW_3$  (*uud*) is  $192 \text{ cm}^{-1}$ , approxi-

mately 2.8 times greater than that in  $W_3$ . Most of the splitting between the single-donor OH stretch modes of  $BW_3$  is the result of the geometrical changes induced by the benzene rather than the coupling between the localized OH stretch modes. In particular, the interaction with the benzene causes large spreads in the single-donor OH and OO distances (Fig. 1), which in turn causes the frequencies of the three single-donor OH stretch local modes to differ appreciably. In order to verify this interpretation, the vibrational frequencies were calculated for each of the three distorted  $W_2$  species that can be generated from the optimized structure for  $BW_3$ . These calculations give single-donor OH stretch frequencies of  $3706$ ,  $3651$ , and  $3561 \text{ cm}^{-1}$ , with an overall splitting of  $145 \text{ cm}^{-1}$ , in fairly good agreement with that calculated (and measured) for the  $BW_3$  cluster.

## 4 Discussion

The results presented in this study demonstrate that the interaction of the  $W_3$  cluster with a benzene molecule causes significant changes in the OH stretch IR spectrum of  $W_3$ . In the single-donor OH stretch regions these changes include an increased splitting between the three transitions and a redistribution of the intensities (with all three lines acquiring nearly equal intensity). The changes in this portion of the spectrum are the result of localization of the normal modes, which, in turn, is due to the structural changes in the  $W_3$  cluster induced by the interaction with the benzene molecule. In the free OH region of the spectrum, the most significant change in going from  $W_3$  to  $BW_3$  is the shift to lower frequency and the large increase in intensity in the mode associated with the OH group  $\pi$  H-bonded to the benzene ring.

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## References

1. Dyke TR, Mack KM, Muentner JS (1977) *J Chem Phys* 66:498; Odutola JA, Dyke TR (1980) *J Chem Phys* 72:5062
2. Fraser GT (1991) *Int Rev Phys Chem* 10:189 and references therein
3. Zwart E, Ter Meulen JJ, Meerts WL (1991) *J Mol Spectrosc* 147:27
4. Coudert LH, Lovas FJ, Suenram RD, Hougen JT (1987) *J Chem Phys* 87:6290
5. Huang ZS, Miller RE (1989) *J Chem Phys* 91:6613
6. Huisken F, Kaloudis M, Kulcke A (1996) *J Chem Phys* 104:17
7. Busarow KL, Cohen RC, Blake GA, Laughlin KB, Lee YT, Saykally RJ (1989) *J Chem Phys* 90:3937
8. Pugliano N, Cruzan JD, Loeser JG, Saykally RJ (1993) *J Chem Phys* 98:6600
9. Pugliano N, Saykally RJ (1992) *Science* 257:1937; Liu K, Elrod MJ, Loeser JG, Cruzan JD, Pugliano N, Brown MG, Rzepiela J, Saykally RJ (1994) *Faraday Discuss* 97:35; Liu K, Loeser JG, Elrod MJ, Host BC, Rzepiela JA, Pugliano N, Saykally RJ (1994) *J Am Chem Soc* 116:3507

10. Liu K, Cruzan JD, Saykally RJ (1996) *Science* 271:929; Cruzan JD, Braly LB, Liu K, Brown MG, Loeser JG, Saykally RJ (1996) *Science* 271:59
11. Xantheas SS, Dunning TH Jr (1993) *J Chem Phys* 99:8774; Xantheas SS (1995) *J Chem Phys* 102:4505; Xantheas SS (1996) *J Chem Phys* 104:8221
12. Fowler JE, Schaeffer HF III (1995) *J Am Chem Soc* 117:446
13. Mhin BJ, Kim J, Lee S, Lee JY, Kim KS (1994) *J Chem Phys* 100:4484; Kim J, Mhin BJ, Lee SJ, Kim KS (1994) *Chem Phys Lett* 219:243
14. van Duijneveldt-van de Rijdt JGCM, van Duijneveldt FB (1993) *Chem Phys* 175:271
15. Feller D (1992) *J Chem Phys* 96:6104
16. Kim J, Lee JY, Lee S, Mhin BJ, Kim KS (1995) *J Chem Phys* 102:310
17. Chakrovorty SJ, Davidson ER (1993) *J Phys Chem* 97:6373
18. van Duijneveldt-van de Rijdt JGCM, van Duijneveldt FB (1992) *J Chem Phys* 97:5019
19. Tsai CJ, Jordan KD (1993) *Chem. Phys Lett* 213:181
20. Kim K, Jordan KD (1994) *J Phys Chem* 98:10089
21. Bürgi T, Graf S, Leutwyler S, Klopper W (1995) *J Chem Phys* 103:1077; Klopper W, Schütz M, Lüthi H, Leutwyler S (1995) *J Chem Phys* 103:1085
22. Gregory JK, Clary DC (1997) *J Phys Chem* (in press)
23. Gotch AJ, Zwier TS (1992) *J Chem Phys* 96:3388; Garrett AW, Zwier TS (1992) *J Chem Phys* 96:3402
24. Ergdah A, Nelander B (1987) *J Chem Phys* 91:2253
25. Wanna AJ, Menapace JA, Bernstein ER (1986) *J Chem Phys* 85:1795
26. Bornsen KO, Selzle HL, Schlag EW (1990) *Z Naturforsch A* 45:1217
27. Gotowsky HS, Emilsson T, Arunan E (1993) *J Chem Phys* 99:4883
28. Suzuki S, Green PG, Bumgarner RE, Dasgupta S, Goddard WA III, Blake GA (1992) *Science* 256:942
29. Maxton PM, Schaeffer MW, Felker PM (1995) *Chem Phys Lett* 241:603
30. Pribble RN, Zwier TS (1994) *Science* 265:75
31. Pribble RN, Zwier TS (1994) *Faraday Discuss* 97:229
32. Pribble RN, Garret AW, Haber K, Zwier TS (1995) *J Chem Phys* 103:531
33. Augspurger JD, Dykstra CE, Zwier TS (1992) *J Phys Chem* 96:7252; *ibid* (1993) 97:980
34. Kim K, Jordan KD, Zwier TS (1994) *J Am Chem Soc* 116:11568
35. Fredericks SY, Jordan KD, Zwier TS (1996) *J Phys Chem* 100:7810
36. Sorenson JM, Gregory JK, Clary DC (1997) *J Chem Phys* 106:849
37. Curtiss LA, Frurip DJ, Blander M (1979) *J Chem Phys* 71:2703
38. Cheng BM, Grover JR, Walters EA (1995) *Chem Phys Lett* 232:364
39. Wales DJ (1993) *J Am Chem Soc* 115:11180
40. Wales DJ (1997) *Encyclopedia of computational chemistry* (in press)
41. Schütz M, Bürgi T, Leutwyler S, Bürgi HB (1993) *J Chem Phys* 99:5228